

# Scintillator detectors with long WLS fibers and multi-pixel photodiodes

O. Mineev,<sup>1,\*</sup> Yu. Kudenko,<sup>1</sup> Yu. Musienko,<sup>1</sup> I. Polyansky,<sup>1,2</sup> and N. Yershov<sup>1</sup>

<sup>1</sup>*Institute for Nuclear Research of RAS*

<sup>2</sup>*Center of Perspective Technology and Apparatus, 107076 Moscow, Russia*

## Abstract

We have studied the possibility of using Geiger mode multi-pixel photodiodes to read out long scintillator bars with a single wavelength-shifting fiber embedded along the bar. This detector configuration can be used in large volume detectors in future long baseline neutrino oscillation experiments. Prototype bars of 0.7 cm thickness and different widths have been produced and tested using two types of multi-pixel photodiodes: MRS APD (CPTA, Moscow) and MPPC (Hamamatsu). A minimum light yield of 7.2 p.e./MeV was obtained for a 4 cm wide bar.

Keywords: Scintillators, WLS fibers, avalanche photodiodes

---

\*Corresponding author. E-mail: oleg@inr.ru

## I. INTRODUCTION

Scintillator bars with wavelength-shifting (WLS) fibers and opto-electronic readout are considered an established technology for massive neutrino tracking calorimeters in long-baseline neutrino oscillation experiments. The MINOS experiment [1] employs extruded bars of  $1 \times 4.1 \times 800 \text{ cm}^3$  size with 9 m long WLS fibers. The SciBar detector [2] in the K2K experiment was built with the same technology as MINOS but using shorter bars of  $1.3 \times 2.5 \times 300 \text{ cm}^3$  size and 3.6 m long WLS fibers of 1.5 mm diameter. A fine-grained detector in the Minerva experiment [3] is made of triangular-shaped 3.5 m long strips and WLS fibers of 1.2 mm diameter. All of these detectors use multi-anode PMTs for optical readout.

The neutrino far detector of the “off-axis superbeam” experiment Nova [4] will be composed of liquid scintillator encased in 15.5 m long rigid PVC extrusion cells. The scintillator cell is readout by a 30 m long U-shaped WLS fiber of 0.7 mm diameter into avalanche photodiodes (APDs), similar to the APDs developed for the CMS detector. These devices have higher quantum efficiency than photomultipliers, but low gain and high intrinsic noise.

A magnetized iron and scintillator sampling calorimeter with a fiducial mass of about 100 kt and a cross sectional area of  $\sim 14 \times 14 \text{ m}^2$  with a 1 T dipole field is considered as a baseline option for a future Neutrino Factory detector. An alternative option is a magnetized, totally active highly segmented scintillator detector of about the same cross sectional area [5]. Both detectors will consist of a large number of readout channels that require the usage of very compact, insensitive to magnetic field photosensors with a high efficiency to the green light emitted from WLS fibers. Multi-pixel Geiger mode avalanche photodiodes are considered as a possible optical readout in these detectors. Detailed information about such devices and their basic principle of operation can be found in ref. [6]. The first application of such photosensors in a large experiment has been done in the near neutrino detector [7] of the long baseline experiment T2K [8] where approximately 56000 Multipixel Photon Counters (MPPCs) [9] are used.

The real challenge lies in the required fine granularity and size of the detectors in these new experiments. Each individual element should provide the ability to detect minimum ionizing particles with high efficiency in such a large detector system. In this paper, results of measurements using scintillator detectors read out with long WLS fibers and Geiger mode

multi-pixel avalanche photodiodes are presented.

## II. PHOTSENSOR SPECTRAL SENSITIVITY

Good spectral matching of the light signal and a photosensor quantum efficiency is one of the most important factors needed to obtain a high light yield in scintillator detectors. We measured the spectral transmittance of Y11 Kuraray WLS fibers as well as the spectral sensitivities of two types of multi-pixel Geiger photodiodes: Hamamatsu MPPC [9] and MRS APD [10].

A 796-pixel MRS APD (type CPTA 151-30) is produced by the CPTA company (Moscow, Russia). The sensitive area of this device is approximated by a circle of diameter 1.28 mm. The pixel size is  $43 \times 43 \mu\text{m}^2$ . A typical gain is close to  $10^6$ , the combined crosstalk and afterpulse probability is estimated to be about 10%, and the dark rate is 1.2-1.5 MHz. A customized 667-pixel MPPC (type S10362-13-050C) with an active area of  $1.3 \times 1.3 \text{ mm}^2$  was developed by Hamamatsu for the near neutrino detector of the T2K experiment. The pixel size of the MPPC is  $50 \times 50 \mu\text{m}^2$ . It should be noted that the response of multi-pixel Geiger photodiodes depends on the overvoltage  $\Delta V$ , i.e. excess of bias voltage over the avalanche breakdown value. MPPCs biased at a typical overvoltage of  $\Delta V = 1.2 \text{ V}$  at  $T=25^\circ\text{C}$  are characterized by the following parameters: a gain of  $7 \times 10^5$ ; an average dark rate of 700 kHz; measured crosstalk and afterpulse probabilities of 9-12% and 14-16% respectively, with a combined value of 20-25%. A front view of both photosensors is shown in Fig. 1.

A spectrophotometer calibrated with a PIN-diode was used to measure the photon detection efficiency (PDE) of both MRS APD and MPPC [11]. The PDE spectra were measured at  $\Delta V=1.2 \text{ V}$  for the MPPC and  $\Delta V=2 \text{ V}$  for the MRS APD. The spectrophotometer light intensity was reduced until the maximum photosensor current was only  $\sim 30\%$  greater than the dark current of a photodiode to avoid nonlinearity effects caused by the limited number of pixels. Comparing the measured current with the calibrated PIN-diode photocurrent we obtain the relative spectral sensitivity. Then the spectral response was normalized using the reference PDE points obtained with LED's at 410 and 515 nm. Average number of photons in LED pulses was measured using a calibrated PMT XP2020. The spectral sensitivity could then be expressed in absolute values without the crosstalk, afterpulse and dark rate

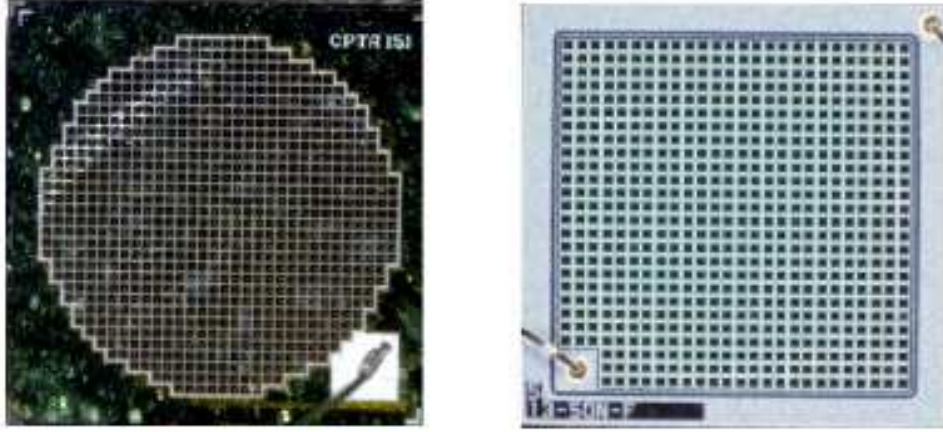


FIG. 1: Front views of MRS APD (left) and MPPC (right).

contributions. The results are shown in Fig. 2. Both MRS APD and the MPPC have similar values of PDE at the Y11 emission peak of 515 nm measured using a 150 cm long Y11 fiber. However, the maximum sensitivity of the MRS APDs is shifted to the red wavelengths around 650 nm while the MPPC sensitivity peaks in the blue region around 450 nm.

The different spectral sensitivities of the MRS APD and MPPC within the visible light range can be explained by their different semiconductor layer structures. The MRS APD uses a  $n^+-p-p^+$  structure while the simplified scheme from the Hamamatsu MPPC catalogue shows the  $p^+-p^--n^{++}$  structure which is usually associated with a peak spectral sensitivity for blue photons.

### III. LONG FIBER STUDY

#### A. Y11 fiber transmittance

The transmittance of 1 mm diameter multi-clad Kuraray WLS Y11(200) S-type fibers [12] was measured for wavelengths of 340-800 nm using clear fibers of 1 mm diameter as a reference. The fiber ends were glued in ferrules and polished. Monochromatic light at a selected wavelength was injected into the test fiber through one end. The intensity of the transmitted light was then measured by a calibrated photodiode attached to the opposite end.

The transmittances of Y11 fibers normalized to the transmittance of a 10 cm reference

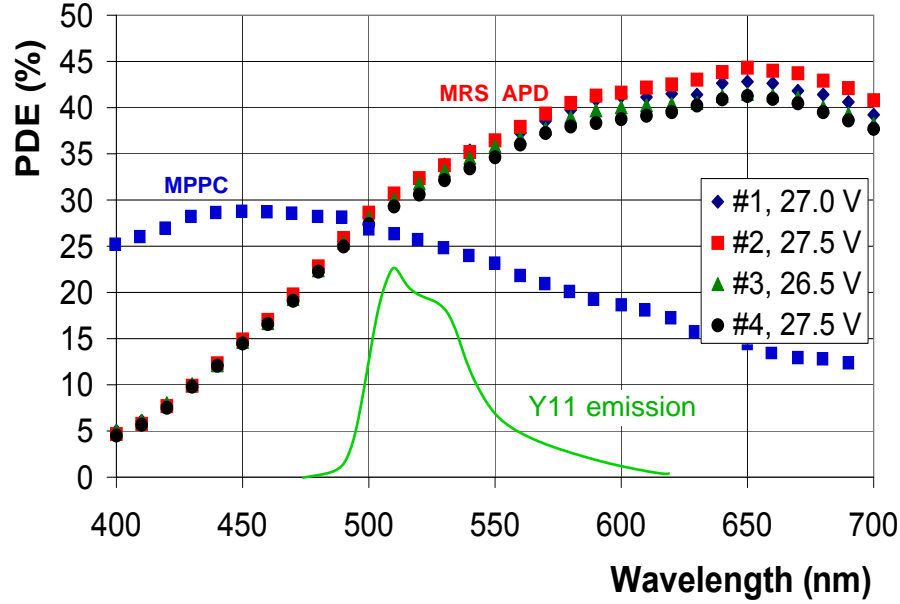


FIG. 2: Photon detection efficiency as a function of wavelength for 4 samples of MRS APDs and an MPPC at 25°C. The bias voltage for each MRS APD was set to an overvoltage of  $\Delta V=2$  V for all tested samples. MPPC was measured at  $\Delta V=1.2$  V. Also shown is the Y11(150) Kuraray fiber emission spectrum (in arbitrary units) for a fiber length of 150 cm (Kuraray specification).

clear fiber are shown in Fig. 3. Better transmittance of 1 m long Y11 fibers than that of the clear fiber of the same length could be explained by some defects or impurities in the tested clear fibers. Fig. 4 shows the transmittance of long Y11 fibers relative to a 1 m long Y11 fiber in the 500-800 nm range. A 15 m long WLS fiber absorbs essentially all the light with wavelengths below 500 nm. For longer wavelengths, the transmittance increases to a maximum level of about 10% at 550 nm. Since the Y11 emission spectrum has wavelengths of 500-550 nm with a low intensity tail extended to over 600 nm (see Fig. 2), one can expect that attenuation of the Y11 signal will be about 95% after the light travels a distance of 15 m.

The dependence of the attenuation length of the Y11 fiber on the wavelength was obtained from the transmittance measurements using a single exponential fit function. Fig. 5 shows the Y11 attenuation length for wavelengths larger than 500 nm. For shorter wavelengths, the light signal was too small in 10 m and 15 m long fibers to obtain reliable values.

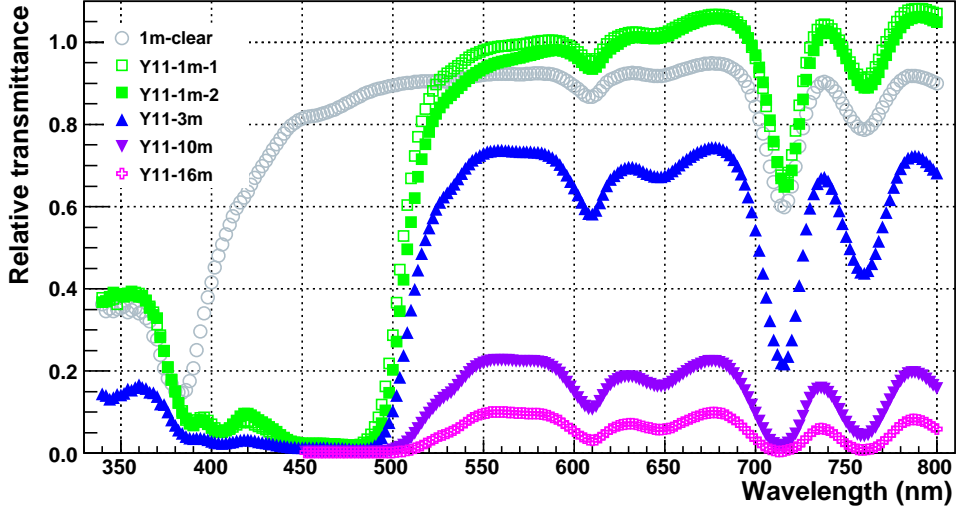


FIG. 3: Relative spectral transmittance of 1 m (two samples), 3 m, 10 m and 15 m long Y11 Kuraray fibers. Transmittance is equal to 1.0 at all wavelengths for a 10 cm long reference clear fiber.

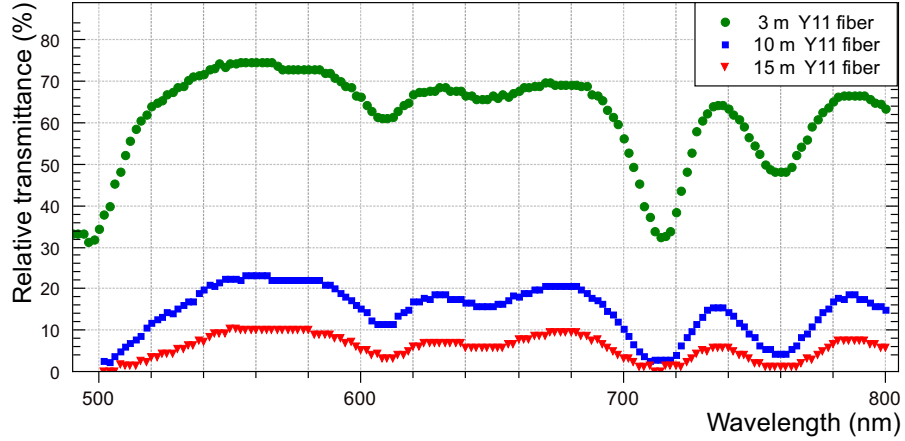


FIG. 4: Relative spectral transmittance for WLS Y11 Kuraray fibers of different lengths. Transmittance is equal to 100% for a 1 m long reference Y11 fiber.

### B. Y11 fiber attenuation

To study the attenuation of the re-emitted light in a 16 m long Y11 fiber, we carried out measurements of the light yield from cosmic ray muons using a small plastic counter. The fiber was embedded with optical grease into a groove machined on the scintillator surface and both ends of the fiber were readout by multi-pixel photodiodes. Photodiode signals

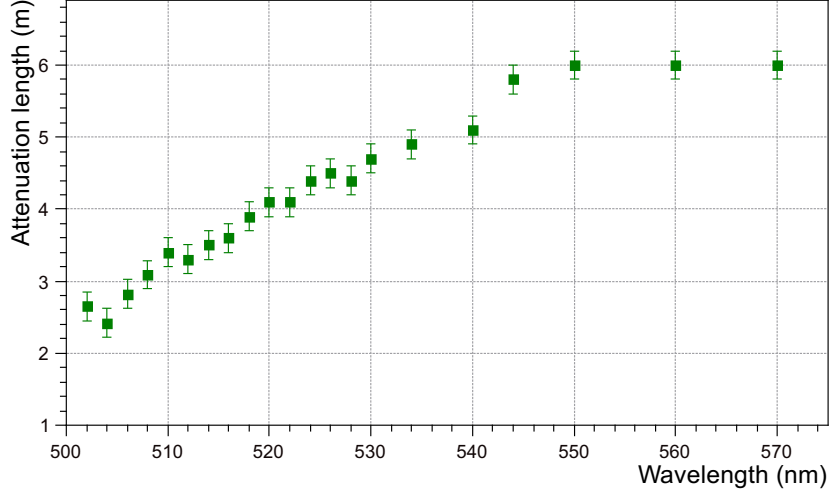


FIG. 5: The attenuation length of Y11 fiber vs wavelength.

were split after the preamplifiers and sent to a charge-integrating LeCroy ADC 2249 and discriminators. The ADC gate was set to 200 ns to ensure that the signal was well inside the integration gate after travelling the 16 m fiber length. The scintillator was moved along the fiber using 1 m steps. The light yield as a function of distance from the photosensors is shown in Fig. 6.

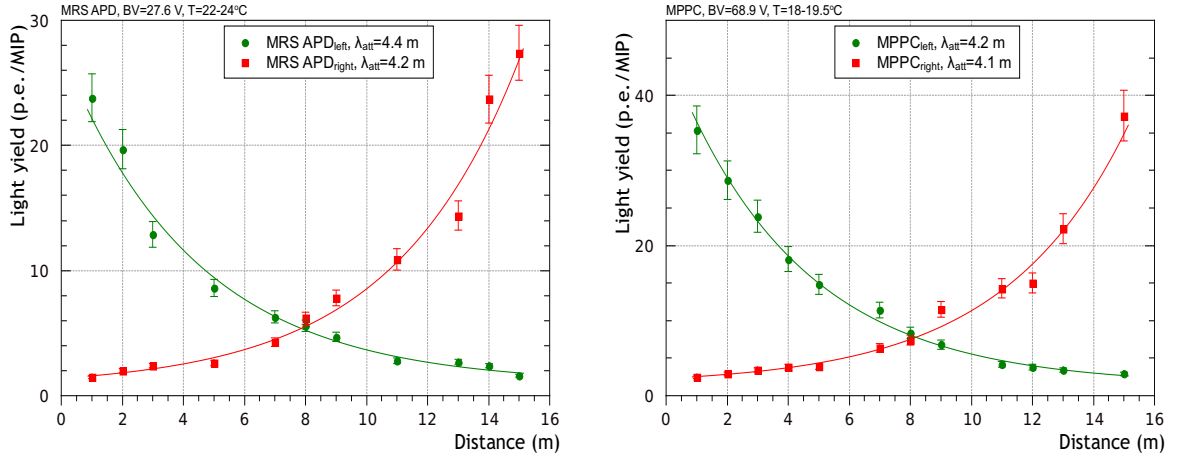


FIG. 6: Light yield for a minimum ionizing particle (MIP) along the 16 m long Y11 fiber where both fiber ends are read out with either MRS APDs (left plot) or MPPCs (right plot). Attenuation is extracted using a single exponential functional fit.

A single exponential function was used to fit for the light yield attenuation along the fiber. An attenuation length of 4.4 m was obtained for the left MRS APD and 4.2 m for the right

one. Attenuation lengths of 4.1 and 4.2 m were obtained for the fiber readout with MPPCs. The shorter attenuation length with MPPCs was expected because the MPPC spectral sensitivity peaks in the blue region. However the difference is rather small and within the accuracy ( $\sigma = 0.4$  m) of the fit. An average attenuation length of 4.2 m corresponds to a light attenuation at an effective wavelength of 520 nm as can be seen in Fig. 5.

To measure the velocity of re-emitted light propagation in a Y11 fiber, a start signal was generated by a scintillator counter from a cosmic ray trigger setup. Another small scintillator tile excited the fiber, and an MRS APD at one fiber end produced the stop signal. The discriminator threshold was set to be fired by 1 photoelectron (p.e.) pulses. The resolution of the LeCroy TDC 2228A was calibrated to be  $50 \pm 0.25$  ps/ch. The time difference between the trigger signal and the signal from the MRS APD is shown in Fig. 7. The velocity of the re-emitted light propagation in the Y11 fiber is found to be  $16.00 \pm 0.08$  cm/ns.

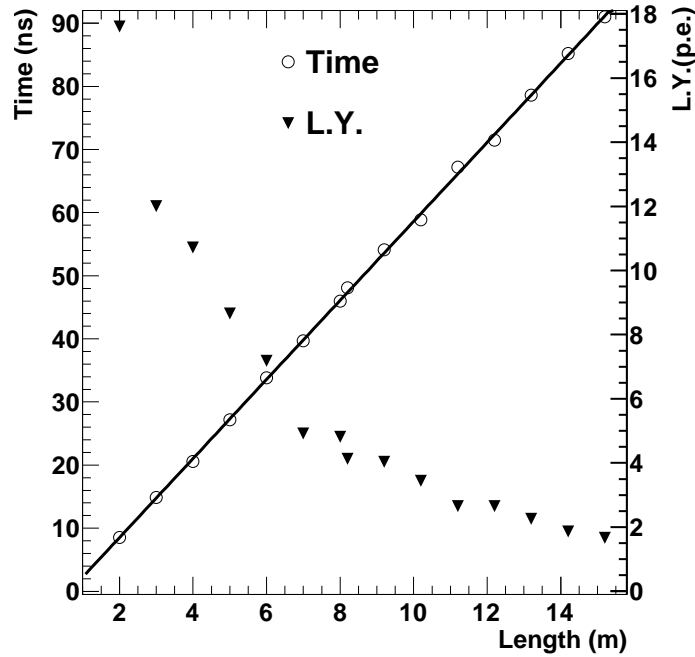


FIG. 7: Time difference between the trigger signal and the signal from the MRS APD along a 16 m long Y11 fiber of 1 mm diameter. The timing error of  $\pm 0.3$  ns is within the radius of the open circles. The light yield in photoelectrons per a minimum ionizing particle is also plotted.



#### IV. MEASUREMENTS OF SCINTILLATOR BARS

Scintillator slabs of  $0.7 \times 20 \times 90$  cm<sup>3</sup> size were extruded at the Uniplast Factory (Vladimir, Russia) and then cut to 90 cm long bars with widths of 1, 2, 3 and 4 cm. The scintillator composition is polystyrene doped with 1.5% of paraterphenyl (PTP) and 0.01% of POPOP. The bars were covered by a chemical reflector by etching the scintillator surface in a chemical agent that results in the formation of a white micropore deposit over the polystyrene [13] surface. The chemical coating is an excellent reflector and in addition it smooths out the rough surface acquired during the cutting process. A 2 mm deep, 1.1 mm wide groove was machined along the center of the bar to accomodate a 16 m long Y11 fiber. Since the tested bar is moved along the fiber, optical grease was used as the coupling between the fiber and the bar.

The test bench for detector measurements is shown in Fig 8. A 16 m long fiber was coiled

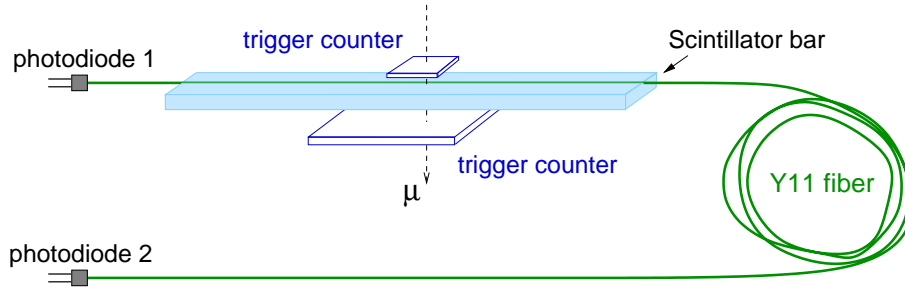


FIG. 8: Test bench for study of the detector read out using a 16 m long Y11 fiber.

inside a light tight box with a minimum bending radius of about 30 cm. The ionization area within the bars was localized to a  $2 \times 2$  cm<sup>2</sup> spot defined by the trigger counter size. The position of the trigger counter was fixed to the center of the 90 cm long bar to optimize scintillation light collection by the fiber. Since a measurement of the 1 cm thick extruded slab yielded an effective attenuation length of the scintillation light in the scintillator of approximately 8.1 cm [14], we estimate that the full scintillation light collection by a WLS fiber occurs within  $\pm 25$  cm from the cosmic ray ionization point.

The light yield along the fiber (sum of both ends) for different bars in photoelectrons (p.e.) per a minimum ionizing particle (MIP) is plotted in Fig. 9. The readout was performed using MRS APDs and the light yield values have been corrected for dark pulses, optical crosstalk and afterpulsing in the photodiodes. The 1 cm wide bar produces the highest light yield

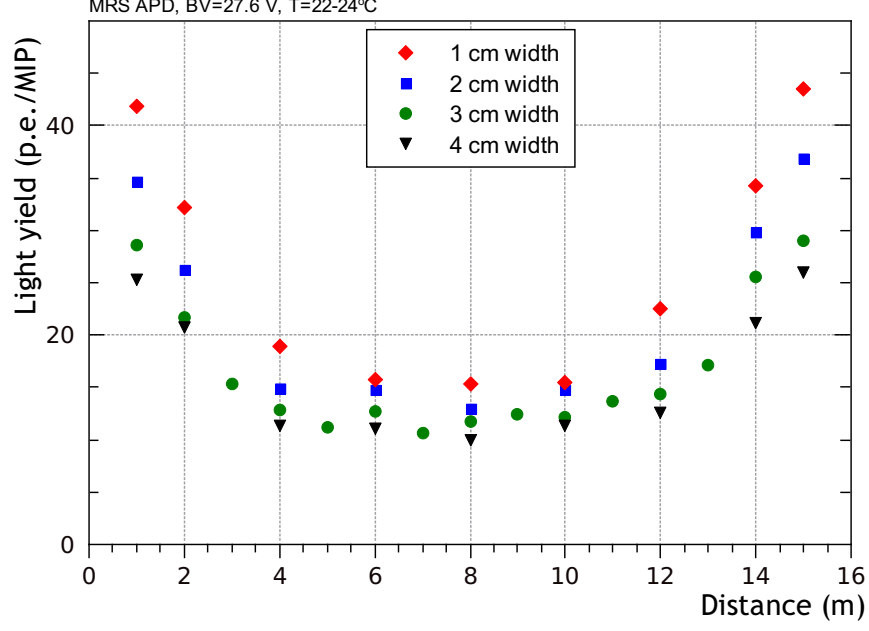


FIG. 9: Total light yield from both fiber ends vs the position along the Y11 fiber for 0.7 cm thick bars of different widths. Measurement errors are estimated to be about 7%.

when compared to the wider bars. The light yield of 15 p.e./MIP in a central point along the 1 cm wide bar at a distance of 8 m from both photosensors drops to 10 p.e./MIP for the 4 cm wide bar. The MIP energy deposit in a bar is 1.4 MeV and corresponds to a minimum light yield of about 7 p.e./MeV in a 0.7 cm thick bar of 4 cm width. Readout using MRS APDs and MPPCs produced similar results within the measurement errors as shown in Fig. 10.

Time resolution in the central part of the fiber was measured to be  $\sigma_t \sim 2.0\text{--}2.2$  ns for the bars with widths ranging from 1 to 4 cm. We made no time walk correction because the discriminator is fired at the arrival of a first photoelectron pulse. Timing is essentially dominated by the Y11 fiber emission decay time. The decay time of Y11 was measured at the center of a 3.5 m long fiber by fitting the time spectrum of a single photoelectron detected at one fiber end while imposing the condition that a few p.e.s were detected at the opposite end. The decay time of Y11 fiber was found to be  $12 \pm 0.5$  ns. Using the light propagation velocity in the fiber the spatial resolution along the bar is estimated to be  $\sigma_x \sim 32\text{--}35$  cm.

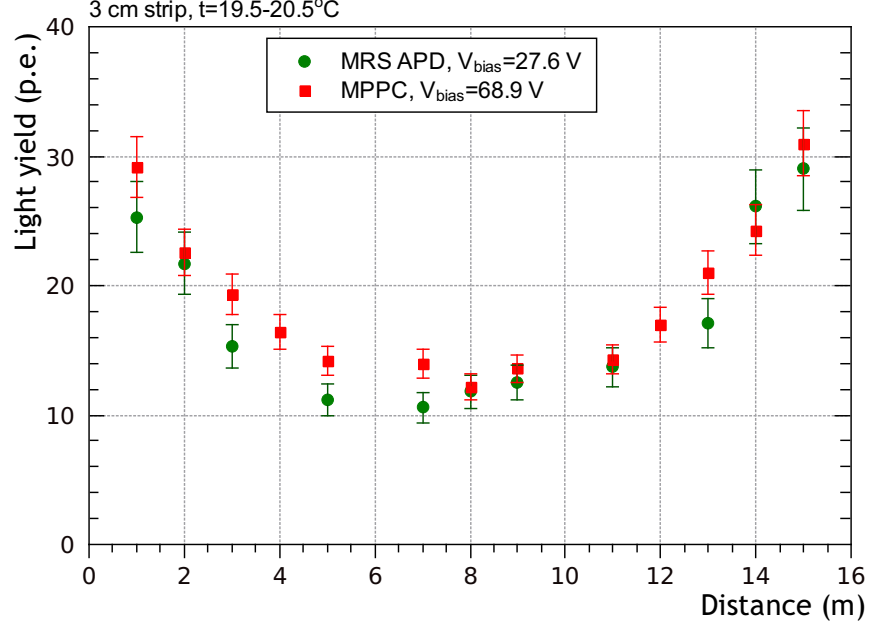


FIG. 10: Light yield (sum of both ends) along the Y11 fiber for the 3 cm wide bar read out with MRS APDs and MPPCs.

## V. CONCLUSION

Scintillator bar readout using a long WLS fiber and multi-pixel photodiodes was studied for a possible usage in large size neutrino detectors. The relative light transmittance for the fiber was measured and compared to the spectral sensitivity of the multi-pixel photodiodes. An attenuation length of 4.2 m at  $\sim 520$  nm enables us to obtain a good signal at a distance of 8 m from an ionization point.

Extruded scintillator bars were read out using a 16 m long Y11 fiber and multi-pixel photodiodes (MRS APDs and MPPCs). The minimum light yield from both fiber ends was measured to be  $\sim 10$  p.e./MIP for the 4 cm wide bar and increased to 15 p.e./MIP for the 1 cm wide bar of the same thickness of 0.7 cm.

The tests have demonstrated that reading out both ends of 16 m long scintillator bars with a single Y11 WLS fiber and multi-pixel photodiodes can provide a high detection efficiency for minimum ionizing particles.

## VI. ACKNOWLEDGEMENTS

This work was supported in part by the “Neutrino Physics” Program of the Russian Academy of Sciences, by the RFBR (Russia)/JSPS (Japan) grant 11-02-92106 and by Science School grant 65038.2010.2.

- 
- [1] D. G. Michael et al, Nucl. Instr. Meth. A596 (2008) 190.
  - [2] K. Nitta et al., Nucl. Instr. Meth. A535 (2004) 147.
  - [3] D. Drakoulakos et al., hep-ex/0405002; K.S. McFarland, Nucl. Phys. Proc. Suppl. 159 (2006) 10.
  - [4] D. Ayres et al., hep-ex/0210005; hep-ex/0503053.
  - [5] T. Abe et al., JINST, 4 (2009) T05001; arXiv:0712.4129.
  - [6] D. Renker and E. Lorenz, JINST 4 (2009) P04004.
  - [7] Yu. Kudenko, Nucl. Instr. Meth. A598 (2009) 289.
  - [8] T2K collaboration, K. Abe et al., arXiv:1106.1238v2 [physics.ins-det].
  - [9] MPPC (Multi-Pixel Photon Counter) is the trademark of HAMAMATSU PHOTONICS K. K. for multi-pixel Geiger mode avalanche photodiodes; M. Yokoyama et al., Nucl. Instr. Meth. A610 (2009) 128; A. Vacheret et al., arXiv:1101.1996 [physics.ins-det].
  - [10] MRS APD is a trademark of the Center of Perspective Technologies and Apparatus (CPTA, Moscow) for multi-pixel photodiodes with a Metal–Resistor–Semiconductor layer structure; <http://www.cpta-apd.ru>
  - [11] Y. Musienko et al., Nucl. Instr. Meth. A567 (2006) 57.
  - [12] Kuraray Co., Methacrylic Resin Division, Tokyo, Japan.
  - [13] Y. Kudenko et al., Nucl. Instr. Meth. A469 (2001) 340.
  - [14] O. Mineev et al., Nucl. Instr. Meth. A577 (2007) 540.